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# Boundary Conditions for the 1-D Non-Classical Transport Equation R. Vasques<sup>1</sup> and K. Krycki<sup>2</sup>

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### 1. Introduction

A non-classical theory has been recently introduced [1] in order to model non-classical particle transport taking place in certain random media. Non-classical transport arises when the probability density function for a particle's distance to collision is not given by an exponential. Assuming that scattering is isotropic, the non-classical transport equation in rod geometry with no internal source is written as

$$\frac{\partial \psi^{\pm}}{\partial s}(x,s) \pm \frac{\partial \psi^{\pm}}{\partial x}(x,s) + \Sigma_t(s)\psi^{\pm}(x,s) = \delta(s)\frac{c}{2}\int_0^\infty \Sigma_t(s')[\psi^+(x,s') + \psi^-(x,s')]ds'.$$
(1)

Here, the path-length s is the distance traveled by the particle since its previous interaction,  $\psi^{\pm}$  represents the angular flux in the directions  $\pm 1$ , c is the scattering ratio, and  $\Sigma_t(s)ds$  represents the probability (ensemble-averaged over all physical realizations) that a particle scattered or born at any point x will experience a collision between  $x \pm s$  and  $x \pm (s + ds)$ .

In this work we investigate for the first time the inclusion of boundary conditions in the non-classical transport equation. This investigation is performed in a homogenized 1-D random periodic system [2] consisting of alternating solid and void layers. Specifically, we address the issue of the assigned value of s in Eq. (1) for particles that enter the system from the exterior. We consider different choices of parameters, providing numerical results for the non-classical equation and comparing these results against "benchmark" numerical results, as well as with the atomic mix and the Levermore-Pomraning models.

### 2. Numerical Results

In order to solve Eq. (1) in the finite  $(0 \le x \le X)$  solid-void 1-D random periodic system given in [2], we use the identity

$$\Sigma_t(s) = \frac{p(s)}{1 - \int_0^s p(s')ds'}$$
(2)

together with the analytically obtained path-length distribution function

$$p(s) = \begin{cases} \frac{\Sigma_{t1}}{\ell} \left[ (2n+1)\ell - s \right] e^{-\Sigma_{t1}(s-n\ell)}, & \text{if } 2n\ell \le s \le (2n+1)\ell \\ \frac{\Sigma_{t1}}{\ell} \left[ s - (2n+1)\ell \right] e^{-\Sigma_{t1}[s-(n+1)\ell]}, & \text{if } (2n+1)\ell \le s \le 2(n+1)\ell \end{cases}$$
(3)

for n = 0, 1, 2, .... Here,  $\Sigma_{t1}$  is the total cross section in the solid material and  $\ell$  is the thickness of each layer in the 1-D system.

For the test problems in this paper we define fixed values of  $\Sigma_{t1}$  and  $\ell$  and present numerical results for different choices of the scattering ratio c. We assume that the incoming fluxes at the boundaries are given by

$$\psi^+(x=0) = 2;$$
 and  $\psi^-(x=X) = 0.$  (4)

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The first approach to model boundary conditions for Eq. (1) is the most intuitively appealing: we simply define that all particles enter the system with the same value of s. The obvious choice in this case is s = 0, as follows:

$$\psi^+(0,s) = 2\delta(s);$$
 and  $\psi^-(X,s) = 0.$  (5)

The second approach assumes that particles enter the system after having already traveled different distances s. One appealing choice consists of using p(s) itself to define the boundary conditions:

$$\psi^+(0,s) = 2p(s);$$
 and  $\psi^-(X,s) = 0.$  (6)

Figure 1 shows numerical results and relative (%) errors obtained for a problem with choice of parameters given by  $\Sigma_{t1} = \ell = 1$ , c = 0.5, and X = 10. We see that the assumption that all particles enter the system with the value s = 0 does not yield an accurate solution, being clearly inferior to the second approach.



Figure 1: Comparison of the ensemble-averaged scalar fluxes and relative (%) errors of the different approaches

The full version of this work will investigate the accuracy of different approaches to model the boundary conditions in this homogenized 1-D random periodic system. We will present solutions for several different choices of parameters, as well as include comparisons with classical models such as Atomic Mix and Levermore-Pomraning.

## References

- E.W. Larsen and R. Vasques, "A Generalized Linear Boltzmann Equation for Non-Classical Particle Transport," J. Quant. Spectrosc. Radiat. Transfer, 112, pp. 619-631 (2011).
- [2] R. Vasques and K. Krycki, "On the Accuracy of the Non-Classical Transport Equation in 1-D Random Periodic Media," Proc. Joint International Conference on Mathematics and Computation, Supercomputing in Nuclear Applications and the Monte Carlo Method, Nashville, TN, April 19-23, on CD-ROM (2015).